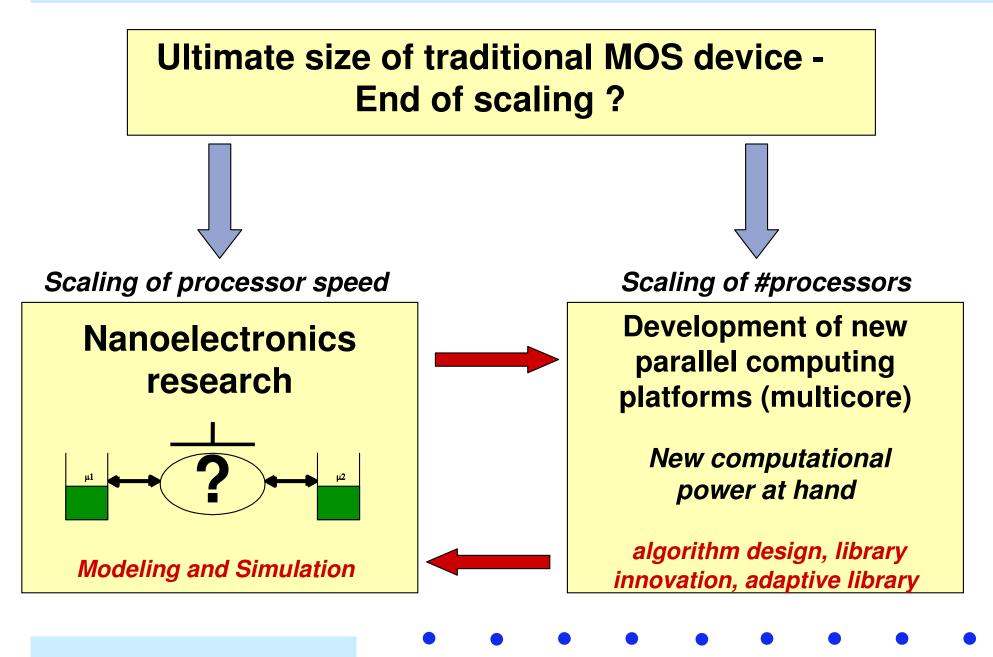
Computational Nanoelectronics

Eric Polizzi

Department of Electrical and Computer Engineering, University of Massachusetts, Amherst, USA

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Computational Nanoelectronics Overview



Computational Nanoelectronics Overview

needed for

- investigating the ultimate size of traditional MOS transistor devices
- leading experiments in the exploration of new class of devices that relies on quantum effects

the need to account for high degrees of detail and realism (quantitative results)

- Multi-scale and multi-dimensional aspect of electron transport and electrostatics effects,
- □ High-fidelity treatment of the contact regions,
- Information of the material at the atomistic level,
- Efficient methodology for the treatment dissipative scattering,

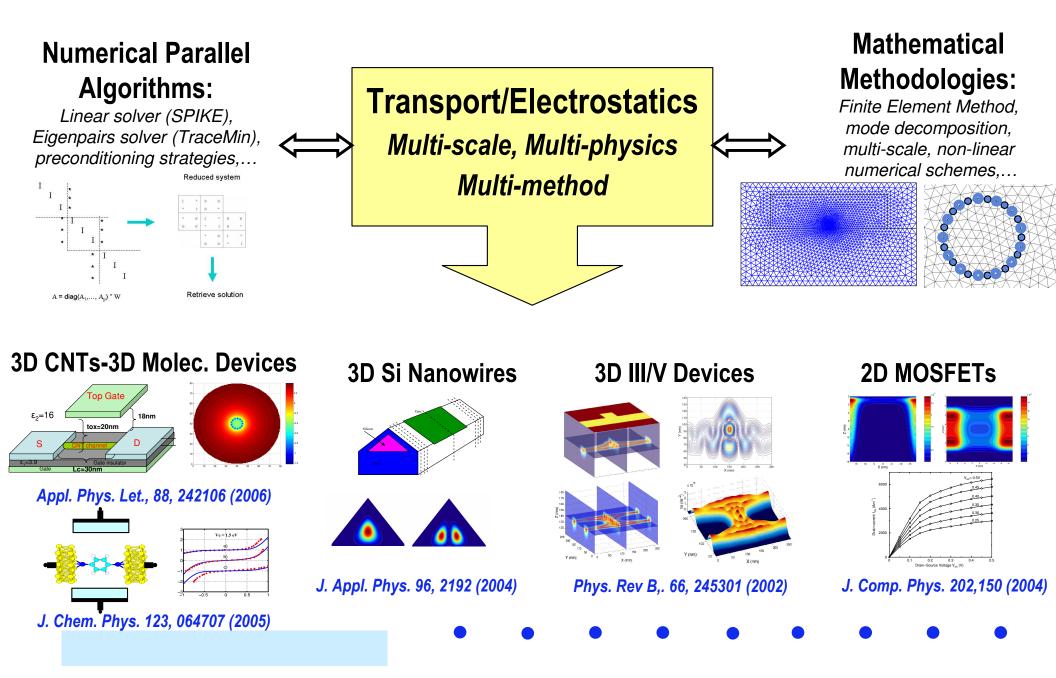
Computationally challenging problems

Computational Nanoelectronics How to proceed ?

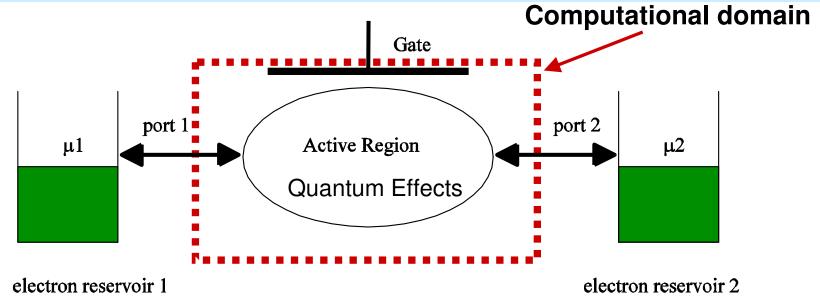
Important features for a Nanoelectronics Simulator:

- Addressing high-fidelity simulations of material & transport phenomena
- Flexible to handle arbitrary devices, different physics, materials, physical dimensions and geometries
 - Prototyping new-class of devices

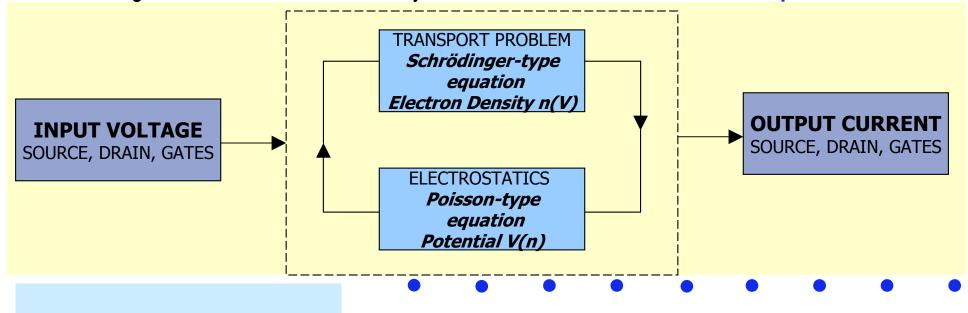
NESSIE (1998-2006) A brief overview



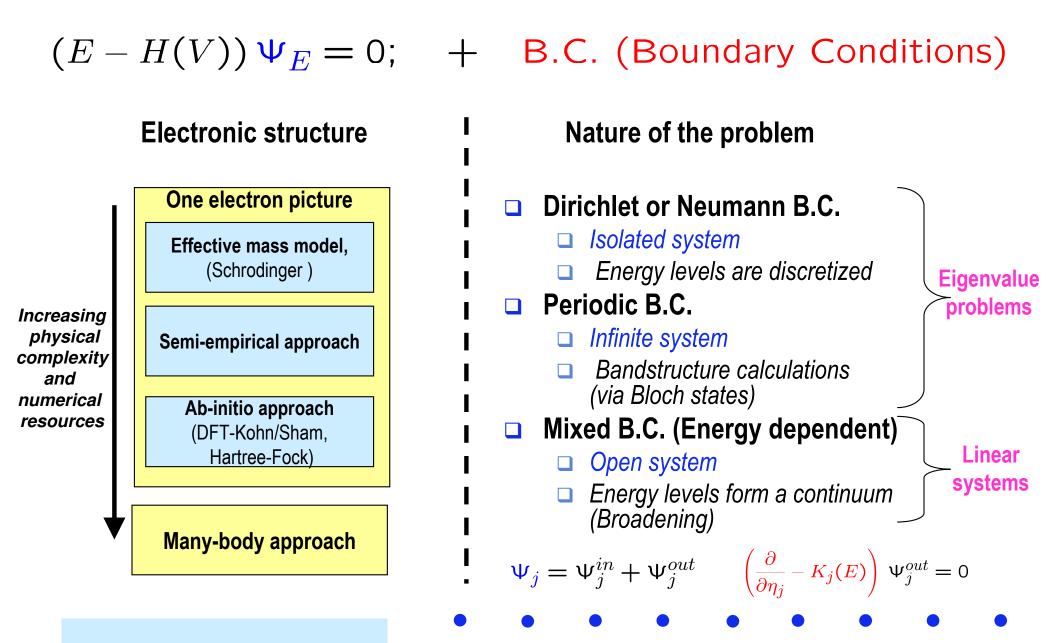
Quantum Transport Modeling: Introduction



Current-Voltage Characteristics obtained by self-consistent simulations: Transport-Electrostatics



Quantum Transport Modeling: Open systems- Ballistic transport



Quantum Transport Modeling: Open systems – FEM discretization

Variational form of the problem (existence and uniqueness: Ben Abdallah et al.):

Find
$$\Psi \in H_g^1 = \{\Psi : \Omega \to C; \Psi \in H^1(\Omega); \Psi = 0 \text{ on } \gamma_0\}$$
, such that $\forall \varphi \in H_g^1 : \frac{\hbar^2}{2} \int_{\Omega} \frac{1}{m^*} \nabla \Psi \nabla \varphi d\Omega + \int_{\Omega} (-qV - E) \Psi \varphi d\Omega - \sum_j \frac{\hbar^2}{2} \int_{\gamma_j} \frac{1}{m^*} \left(K_j \Psi\right) \varphi d\gamma_j = L_{j_0}^{ex}$

Finite Element Discretization: $\Psi(\mathbf{x}) \simeq \sum_{i=1}^{N} \psi_i \omega_i(\mathbf{x})$

$$\left(E[\mathsf{S}] - [\mathsf{H}(\mathsf{V})] - \sum_{\mathsf{j}} [\boldsymbol{\Sigma}_{\mathsf{j}}(\mathsf{E})]\right) \Psi_{j_0}^{ex} = \mathbf{L}_{j_0}^{ex}$$

For the calculation of electron and current densities, it is necessary to solve a large number of those linear systems- Integration over the energy

Numerical Techniques:

Computational Challenges in Nanoelectronics Simulations

Realistic Device Simulations

- Multidimensional geometries
- Detailed description of the electronic structure
- **Dissipative Scattering**
- **Etc...**
- → High numerical costs
- Huge computer resources

For only one point in the I-V curve	Full 2D	Full 3D
Matrix size	O(10 ⁴)	O(10 ⁶)
linear systems to solve by iteration	O(10 ³)	O(10 ³)
# self-consistent iterations	O(10)	O(10)
Simulation time	O(hours)	O(days)

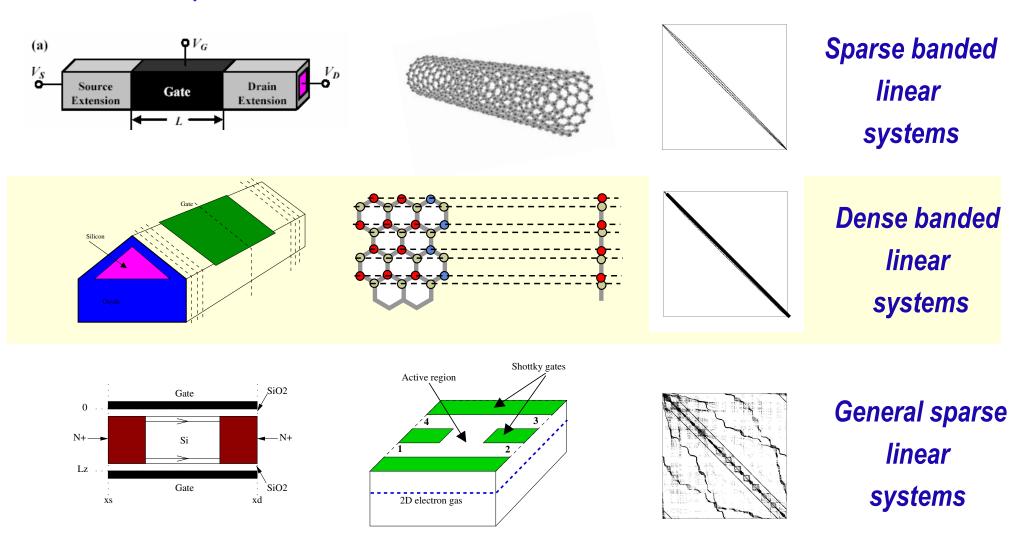
Standard strategy

- Each processor handles many linear systems, one at a time
- Sequential algorithms are inefficient for obtaining I-V curves on large systems

Proposed strategy

- **Development of high**
 - performance numerical parallel algorithms enabling large-scale nanoelectronics simulations such as the linear solver SPIKE

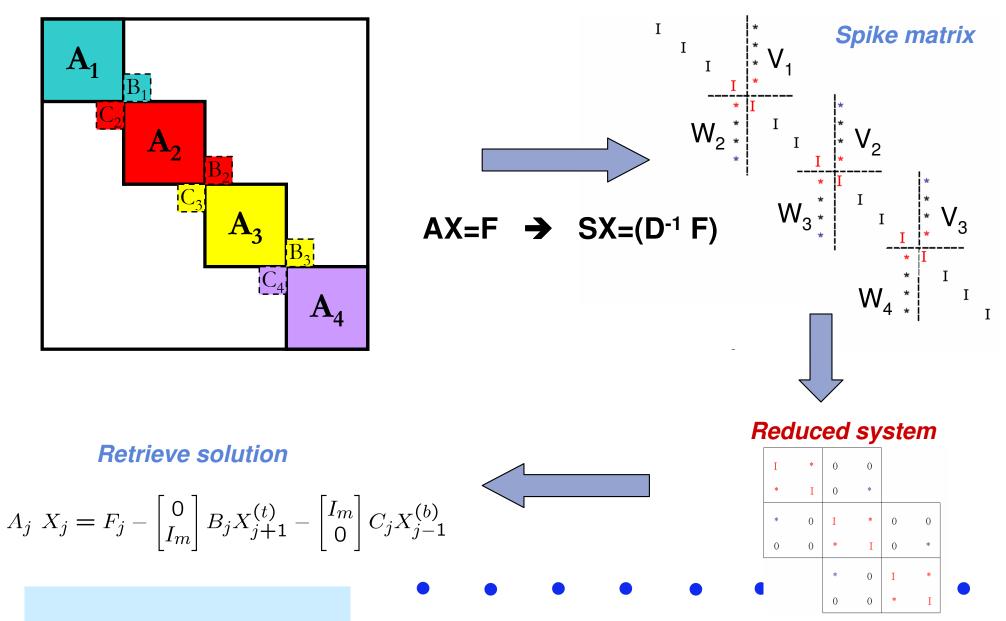
Numerical Techniques: Matrix representation



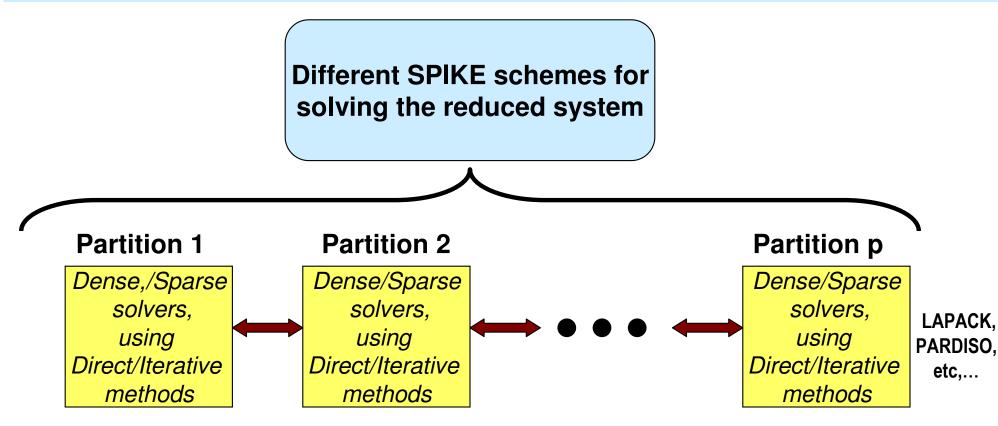
→ Computational Nanoelectronics is in need of efficient parallel banded solvers

SPIKE Overview of the general algorithm

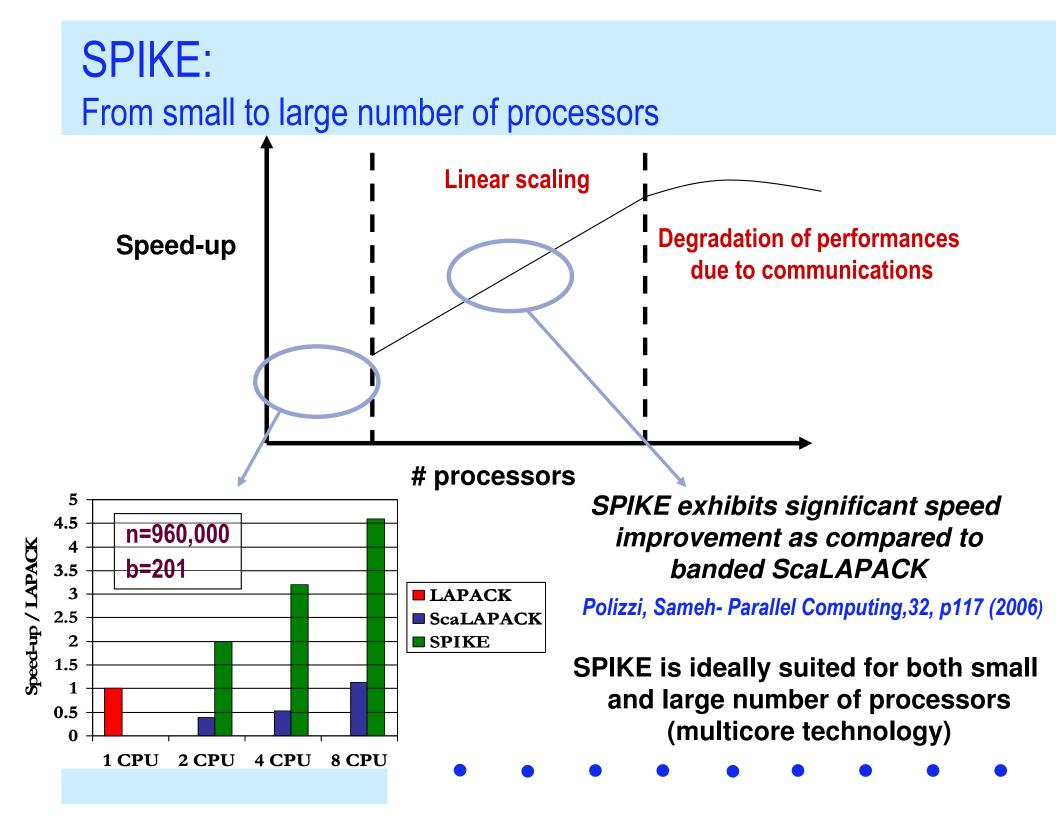
SPIKE is a parallel banded system solver proposed by A. Sameh (77-78), recently revisited.



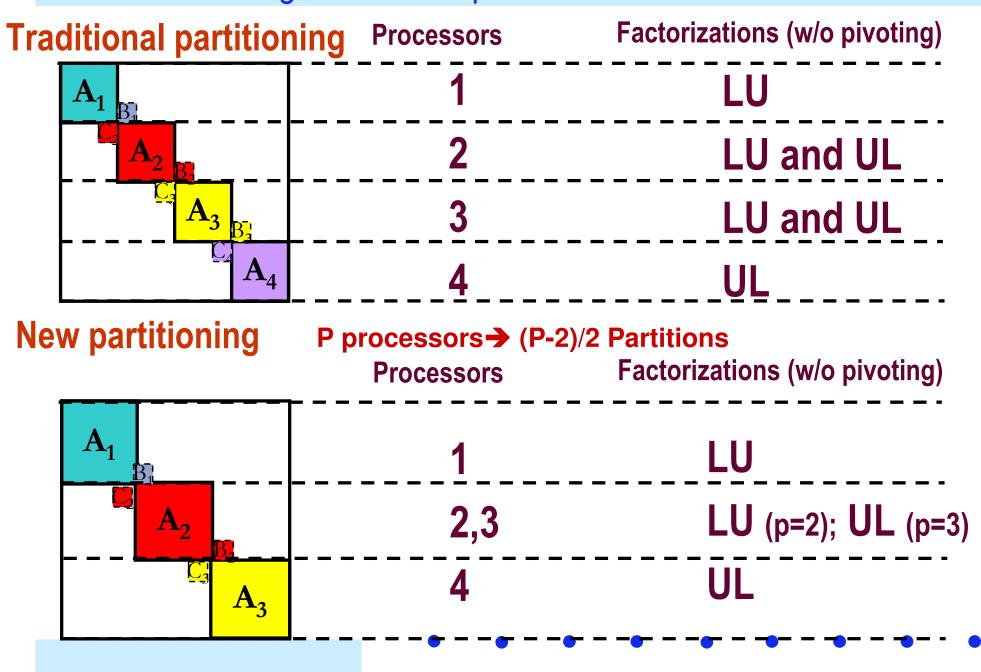
SPIKE: New Implementation - A robust parallel environment



- SPIKE is essentially a family of domain decomposition techniques. The diagonal blocks A_j correspond to independent subdomains, and solving the resulting reduced system corresponds to solving the interface problem,
- **SPIKE** is a hybrid and polyalgorithm (more than 40 different algorithms)



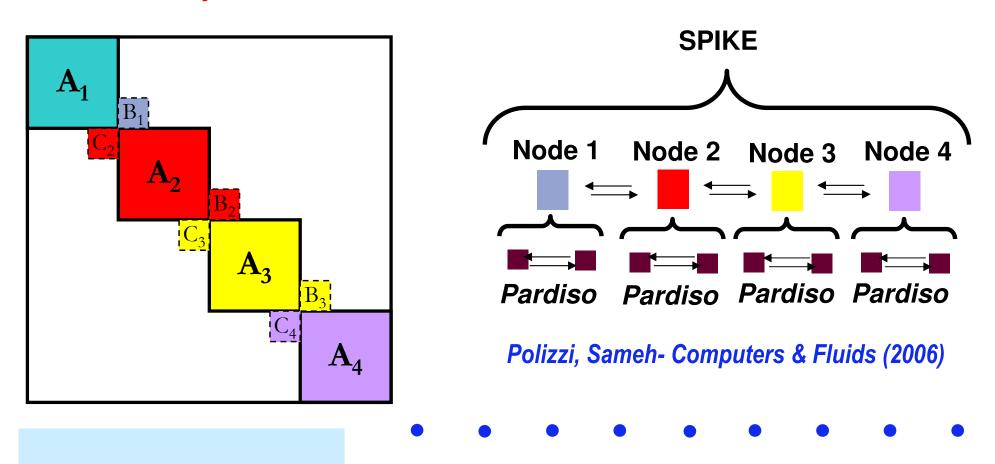
SPIKE: From small to large number of processors



SPIKE:

Banded systems that are sparse within the band

- >We make use of SPIKE "on-the-fly" scheme
- A SPIKE hybrid scheme exhibits better performance than other parallel direct sparse solvers used alone such as Pardiso, SuperLU, MUMPS
 Multilevel of Parallelism can be used to reduce the complexity of the reduced system



SPIKE: Banded systems that are sparse within the band

N=432,000, b= 177, nnz= 7, 955, 116, sparsity of the band: 10.4%

Test (a)	Reord.	Fact.	Solve	Total	Residual
SPIKE 2-nodes	9.48	1.55	4.54	15.57	10 ⁻⁷
SPIKE 4-nodes	5.68	0.77	2.3	8.76	10^{-7}
Trunc. SPIKE (8)	0	0.84	0.14	0.98	10^{-14}

→ MUMPS: time (2-nodes) = 21.35 s; time (4-nodes) = 39.6 s (memory swap)

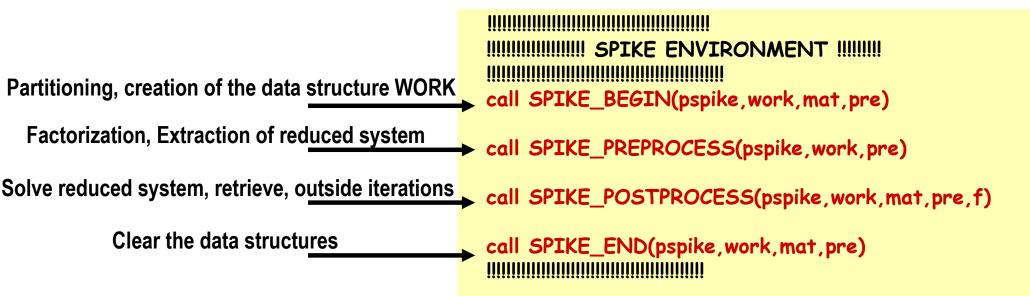
For narrow banded systems, SPIKE will consider the matrix dense within the band.
 Reordering schemes for minimizing the bandwidth can be used if necessary.

N=471,800, b= 1455, nnz= 9, 499, 744, sparsity of the band: 1.4%

Test (b)	Reord.	Fact.	Solve	Total	Residual
SPIKE 2-nodes	18.69	11.06	23.06	52.82	10 ⁻⁷
SPIKE 4-nodes	8.41	6.54	10.85	25.81	10^{-7}

→ Good scalability using "on-the-fly" SPIKE scheme

SPIKE: Current status



- SPIKE v.2 contains ~10K lines of MPI-Fortran90,
- **Documentation (user guide) is ready,**
- More than 40 different schemes for SPIKE available (best choices depend on the properties of the linear systems and/or architectures of the machines).
- **New options: customized partitioning, and customized preconditioner.**
- A SPIKE_TUNE module is currently developed by Intel for automatic selection of the best SPIKE algorithm
- ✤ SPIKE will be soon available on-line

Computational Nanoelectronics Preconditioning strategy

$$(E[S] - [H(\vee)] - \sum_{i} [\Sigma_{j}(E)])X = F, \forall E$$

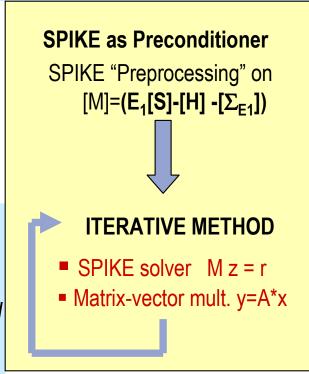
Traditional strategy:

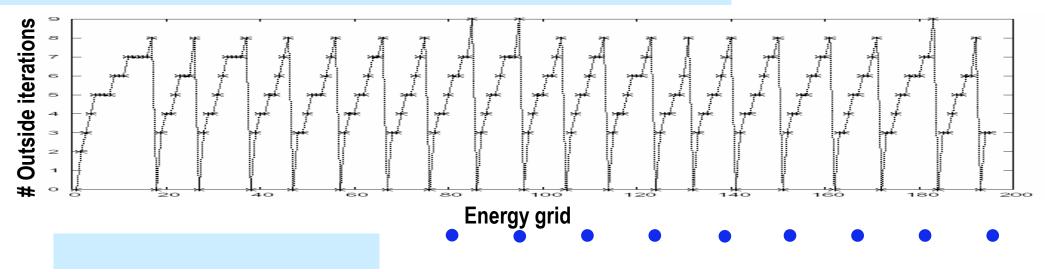
Iterative solvers such as QMR, GMRES, BicgStab, with heuristic preconditioners such as ILUT, SSOR require ~80-100 iterations.

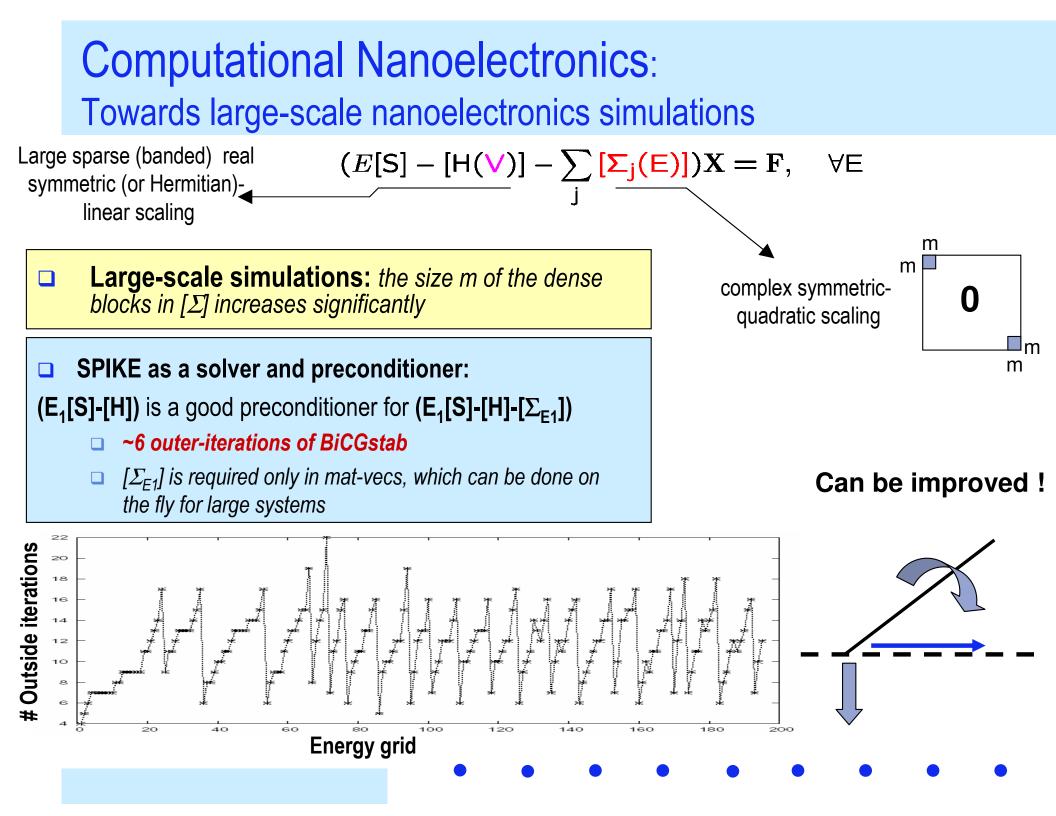
■ We propose to use (E₁[S]-[H] -[Σ_{E1}]) as preconditioner for solving (E₂[S]-[H]-[Σ_{E2}]), if (E2 -E1) < δE ,

the preconditioner is solved in parallel using SPIKE and updated if # of iterations > Nmax

Fast algorithm allowing refinement of the uniform energy grid



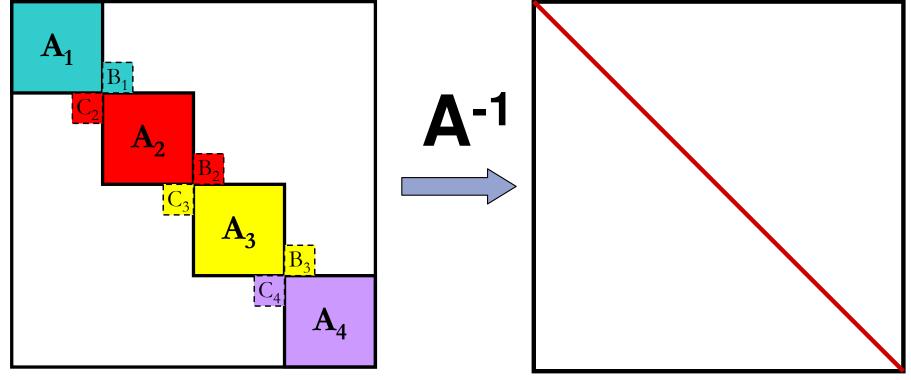




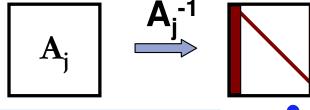
SPIKE: Additional scheme

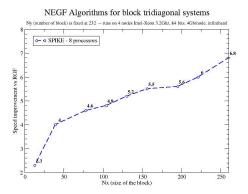
SPIKE can be used for computing the diagonal elements of the inverse

(alternative to standard electronic structure calculations).



 $\hfill\square$ Basic computational module for each ${\rm A_j}$





Computational Nanoelectronics Directions

Modeling Challenges

- Dissipative scattering processes
- Atomistic modeling
- Time dependent

Computational Challenges

- Improvement of the numerical techniques (multiscale, high order discretization schemes, preconditioning strategies...)
- In Need of fast parallel algorithms (SPIKE, etc...)

Applications: Emerging Electronics Devices

- Arbitrary devices, materials and geometries
- CNT and carbon graphene using fully atomistic real-space mesh techniques
- **Full 3D simulations of Nanowire (Si and III-V) and dissipative scattering**

NESSIE as a platform for exploring innovative modeling strategies, the essential numerical methods and high performance algorithms

Acknowledgment: Intel support

Backup

Quantum Transport Modeling: Open systems - derivation of the open B.C.

• CASE 1 (1D device) Frensley (1990)

$$\left. \frac{\partial}{\partial \eta_j} \right|_0 \Psi_j^{out}(\eta_j) = i k_j \Psi_j^{out}(0)$$

• CASE 2 (2D/3D device) Lent, Kirkner (1990)

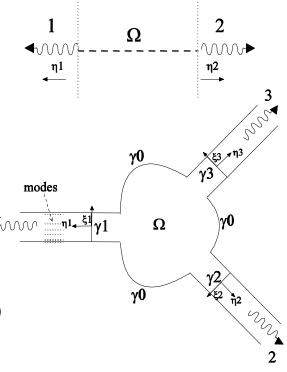
$$\frac{\partial}{\partial \eta_j}\Big|_{0} \Psi_j^{out}(\eta_j,\xi_j) = i \sum_m k_j^m \chi_j^m(\xi_j) \langle \chi_j^m(\xi_j) | \Psi_j^{out}(0,\xi_j) \rangle_{L^2(\gamma_j)}$$

- **GENERAL CASE** (G_j : Green's function of the reservoirs j)
 - Inglesfield (1981)–with Neumann B.C. for G_j on γ_j

$$\frac{\partial}{\partial \eta_j}\Big|_{0} \Psi_j^{out}(\eta_j,\xi_j) = -2 \int_{\gamma_j} \mathbf{G}_j^{-1}(\xi_j;\xi_j') \Psi_j^{out}(0,\xi_j') d\xi_j'$$

– Fisher (1990)– with Dirichlet B.C. for G_j on γ_j

$$\frac{\partial}{\partial \eta_j}\Big|_{0} \Psi_j^{out}(\eta_j,\xi_j) = \frac{1}{2} \int_{\gamma_j} \left(\frac{\partial^2}{\partial \eta_j \partial \eta_j'} G_j(\eta_j,\xi_j;\eta_j'\xi_j') \right) \Psi_j^{out}(0,\xi_j') d\xi_j'$$



Quantum Transport Modeling:

From Jellium to Atomistic models, why using real-space mesh techniques? In contrast to traditional methods such as: empirical tight-binding, LCAO, plane wave scheme

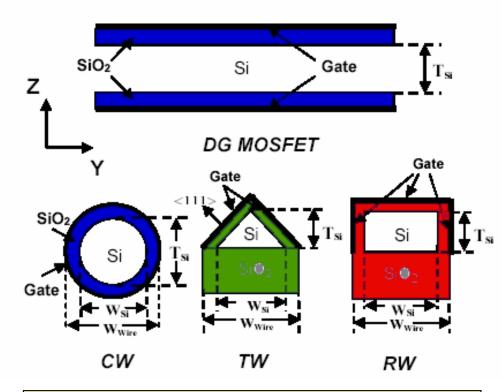
They do not suffer from numerical truncation errors of finite
 Computational expansions when scaled to larger systems.

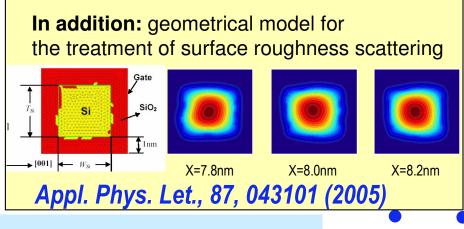
- Materials
 Consistency while solving the Poisson equation for electrostatics
- Science Produce very sparse matrices (banded) and are cast as linear scaling methods.
 - Ease to derive boundary conditions using a mathematical framework

Computational Electronics

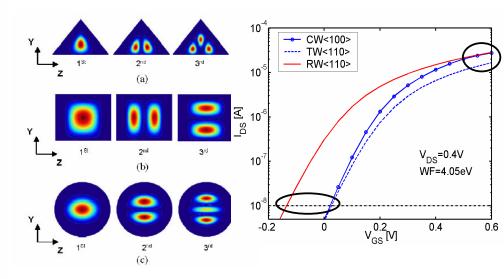
- Ability to consider discontinuities in materials, and the coupling between different physical domains: jellium/atomistic
- FEM is a local basis-oriented approach, with ability to accommodate irregular device geometries, and to consider nonuniform meshes.

Silicon Nanowire Transistors



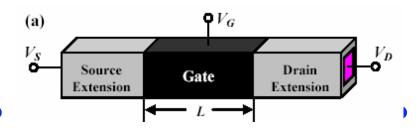


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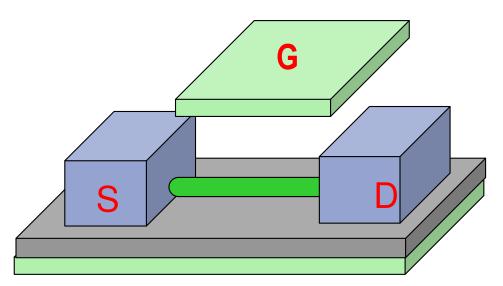


3D Mode space approach

- Asymptotic model (fast uncoupled mode approach or quasi-full dimensional model)
- Full 3D simulation to account for gate leakage in progress



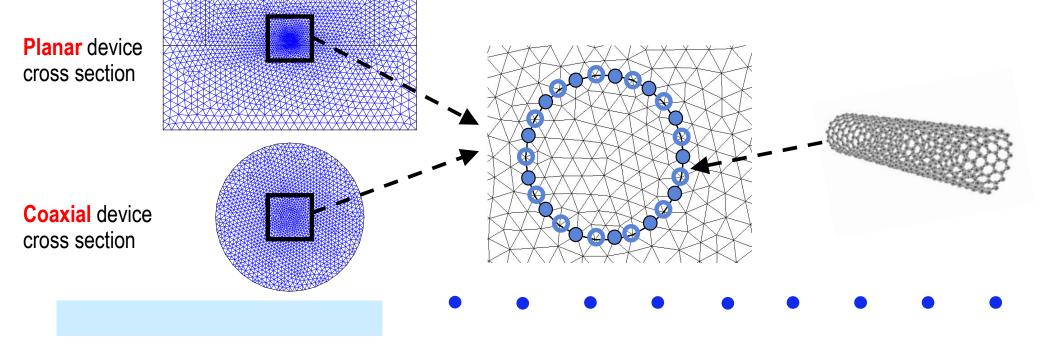
Simulation of Emerging Electronic Devices: Carbon Nanotube Transistors



- □ Huge influence of 3D electrostatics on 1D CNT
- Multiscale-Electrostatics approach for CNTFET:
 - At large scale (transistors), the Carbon atom sites appear as point charges.

Atomistic transport model:

- Currently uses tight-binding p_z orbitals, Mode space or real-space approach; Empirical approach for handling metal-nanotube contacts.
- PDE-based approach under investigation: DFT Kohn/Sham with pseudopotential → hybrid transport model for metal-nanotube contact (jellium-atomistic interface)



Simulation of Emerging Electronic Devices: Carbon Nanotube Transistors

Planar device

